### Handling NIRIS Data: Calibration, Processing, and Inversion

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### **Outline**

- PART I: NIRIS Instrument
	- Brief review of scientific principles and goals
	- Instrumental design
- PART II: NIRIS Data Processing
	- Steps in the NIRIS pipeline
	- Calibration of instrumental crosstalk
- PART III: Inverting NIRIS Data
	- Brief review of Stokes inversions



#### Part I: NIRIS instrument



#### Science enabled by NIRIS

- High-resolution spectroscopy and polarimetry using Zeemansensitive spectral lines in the NIR
	- Zeeman splitting is a function of  $g\lambda^2B$



- Imaging deepest photosphere through base of corona
	- Photosphere opacity minimum at 1.6  $\mu$ m



#### Exploiting the Zeeman effect

- The (orbital) magnetic moments of an atom's electrons interact with an external magnetic field
	- States having different magnetic quantum numbers  $(m_\ell)$  shift in energy in the presence of a magnetic field:  $V = -\vec{\mu} \cdot \vec{B}$
- Resulting Zeeman triplet:

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No information about direction of  $\vec{B}$  from **the intensity spectrum alone**



### Exploiting the Zeeman effect



We can get information about the direction of  $\vec{B}$  by looking at the polarization of the observed spectrum

Stokes parameters:  $I, Q, U, V$ 

- Transverse Zeeman effect
	- $\vec{B}$  perpendicular to LOS
- Longitudinal Zeeman effect
	- $\vec{B}$  along LOS







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#### Measuring magnetic fields with NIRIS

- Method for inferring  $\vec{B}$  using polarization measurements:
	- 1. Measure  $I(\lambda)$ ,  $Q(\lambda)$ ,  $U(\lambda)$ ,  $V(\lambda)$
	- 2. Relate  $I(\lambda)$ ,  $Q(\lambda)$ ,  $U(\lambda)$ ,  $V(\lambda)$  to  $\vec{B}$  using a radiative transfer theory
		- This will include Zeeman effect, Hanle effect (depolarization through scattering), assumptions about solar atmosphere, etc.
	- 3. Find  $\vec{B}$  by solving the inverse problem ("Stokes inversion")
- NIRIS: dual Fabry-Perot system that performs dual-beam polarimetry
	- 2 FPIs are used as a tunable narrow-band filter, scans over spectral line
	- Dual-beam design reduces seeing-induced polarization crosstalk
		- Beam splitter separates orthogonal linear polarizations
		- Combinations of the two resulting images produce  $I, Q, U, V$  images



#### Part II: NIRIS data processing



modulation angle

### Separating  $I, Q, U, V$  images



- Modulator waveplate (frequency  $\omega$ ) modulates polarization states differently
	- Q and U (linear polarization) modulated at  $4\omega$
	- V (circular polarization) modulated at  $2\omega$

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• Observed modulated profile is a combination of I,  $Q$ ,  $U$ ,  $V$ :

$$
I'' = \frac{1}{2} \left[ I' + \frac{Q'}{2} \left( (1 + \cos \delta) + (1 - \cos \delta) \cos 4\theta \right) + \frac{U'}{2} \left( 1 - \cos \delta \right) \sin 4\theta - V' \sin \delta \sin 2\theta \right]
$$
  
waveplate retardation

#### Separating  $I, Q, U, V$  images





• I, Q, U, V signals can be recovered via different linear combinations of integrations  $1 - 8$ 



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### Separating  $I, Q, U, V$  images

- Remember: our camera frame rate isn't exactly matched with the modulator frequency – there will be some offset  $\phi$  when integrating over the frames
	- For a perfectly matched camera frame rate:

$$
Q' = \int_{0}^{\pi/8} I'' d\theta - \int_{\pi/8}^{1} I'' d\theta + \int_{\pi/4}^{1} I'' d\theta + \int_{\pi/2}^{1} I'' d\theta + \int_{\pi/2}^{1} I'' d\theta - \int_{\pi/8}^{1} I'' d\theta - \int_{\pi/8}^{1} I'' d\theta + \int_{\pi/8}^{1} I'' d\theta + \int_{\pi/8}^{1} I'' d\theta
$$
\n
$$
U' = \int_{0}^{\pi/8} I'' d\theta + \int_{\pi/8}^{1} I'' d\theta - \int_{\pi/8}^{1} I'' d\theta - \int_{\pi/8}^{1} I'' d\theta + \int_{\pi/2}^{1} I'' d\theta + \int_{\pi/2}^{1} I'' d\theta + \int_{\pi/2}^{1} I'' d\theta - \int_{\pi/8}^{1} I'' d\theta - \int_{\pi/8}^{1} I'' d\theta
$$
\n
$$
= \begin{array}{ll}\n\text{Our Case:} \\
\text{Our case:} \\
\alpha^{1/8+\phi} & \pi^{1/4+\phi} \\
Q'' = \int_{\phi}^{1} I'' d\theta - \int_{\pi/8+\phi}^{1} I'' d\theta + \int_{\pi/4+\phi}^{1} I'' d\theta + \int_{\pi/2+\phi}^{1} I'' d\theta + \int_{\pi/2+\phi}^{1} I'' d\theta - \int_{\pi/8+\phi}^{1} I'' d\theta - \int_{\pi/8+\phi}^{1} I'' d\theta + \int_{\pi
$$

 $\blacktriangleright$  Leakage between  $Q$  and  $U$ 

Our case:

$$
Q'' = \int_{\phi}^{\pi/8+\phi} I'' d\theta - \int_{\frac{\pi}{8}+\phi} I'' d\theta - \int_{\frac{\pi}{8}+\phi} I'' d\theta + \int_{\frac{3\pi}{8}+\phi} I'' d\theta + \int_{\frac{3\pi}{8}+\phi} I'' d\theta + \int_{\frac{3\pi}{8}+\phi} I'' d\theta - \int_{\frac{5\pi}{8}+\phi} I'' d\theta - \int_{\frac{5\pi}{8}+\phi} I'' d\theta - \int_{\frac{3\pi}{8}+\phi} I'' d\theta + \int_{\frac{\pi}{8}+\phi} I'' d\theta + \int_{\frac{\pi}{8}+\phi} I'' d\theta
$$
  
=  $Q' \cos 4\phi - U' \sin 4\phi$ 

- Similar effect for  $U$  and  $V$
- Correct for leakage caused by phase offset by multiplying the Stokes vector by a restoration matrix:  $M_{\phi}$

$$
\begin{bmatrix} I' \\ Q' \\ \text{signal that has} \\ \text{passed through} \\ \text{optical chain} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos 4\phi & \sin 4\phi & 0 \\ 0 & -\sin 4\phi & \cos 4\phi & 0 \\ 0 & 0 & 0 & \sec 2\phi \end{bmatrix} \begin{bmatrix} I'' \\ Q'' \\ U'' \\ V'' \end{bmatrix}
$$
\nMeanured at detector optical chain



- NIRIS uses many different optical elements that are difficult to model
- We are trying to recover the "true" Stokes vector  $\vec{S}$ , but the polarization of the light entering the telescope is altered as it passes through the telescope and relay optics
- The way an optical element alters polarization of light is described by a  $4\times4$ Mueller matrix M:

$$
\vec{S}' = \mathbf{M}\vec{S}
$$

- Every element of the NIRIS relay optics alters the polarization of the incident beam
	- We can quantify the total effect of an optical chain using one Mueller matrix  $M$ :

$$
\mathbf{M} = \mathbf{M}_1 \mathbf{M}_2 \mathbf{M}_3 \dots \mathbf{M}_n
$$



• Starting with the telescope Mueller matrix **T** (combined effects of primary and secondary mirrors), the signal after phase correction is:

$$
\vec{S}' = MS = [M_{\text{rest}}R(HA)M_4R(\text{dec})M_3T]\vec{S}
$$

• Trace the path of the light:



• Starting with the telescope Mueller matrix T (combined effects of primary and secondary mirrors), the signal after phase correction is:

$$
\vec{S}' = MS = [M_{\text{rest}}R(HA)M_4R(\text{dec})M_3T]\vec{S}
$$

- Trace the path of the light:
	- 1.  $\vec{S}$  is altered as the beam passes through the primary and secondary mirrors, whose effects are quantified by T



• Starting with the telescope Mueller matrix T (combined effects of primary and secondary mirrors), the signal after phase correction is:

$$
\vec{S}' = MS = [M_{\text{rest}}R(\text{HA})M_{4}R(\text{dec})M_{3}T]\ \vec{S}
$$

- Trace the path of the light:
	- 1.  $\vec{S}$  is altered as the beam passes through the primary and secondary mirrors, whose effects are quantified by T
	- 2. The beam reflects off mirror 3  $(M<sub>3</sub>)$  and is sent down the declination axis of the telescope, resulting in a rotation in dec,  $R(dec)$



• Starting with the telescope Mueller matrix T (combined effects of primary and secondary mirrors), the signal after phase correction is:

 $\vec{S}' = \mathbf{MS} = [\mathbf{M_{rest}R(HA)M_4R(dec)M_3T}] \ \vec{S}$ 

- Trace the path of the light:
	- 1.  $\vec{S}$  is altered as the beam passes through the primary and secondary mirrors, whose effects are quantified by T
	- 2. The beam reflects off mirror 3  $(M<sub>3</sub>)$  and is sent down the declination axis of the telescope, resulting in a rotation in dec,  $R(dec)$
	- 3. The beam reflects off mirror 4 ( $M<sub>4</sub>$ ) and is sent down the polar axis of the telescope, resulting in a rotation in hour angle,  $R(HA)$



• Starting with the telescope Mueller matrix T (combined effects of primary and secondary mirrors), the signal after phase correction is:

 $\vec{S}' = \mathbf{MS} = [\mathbf{M_{rest}R}(\text{HA})\mathbf{M_4R}(\text{dec})\mathbf{M_3T}] \ \vec{S}$ 

- Trace the path of the light:
	- 1.  $\vec{S}$  is altered as the beam passes through the primary and secondary mirrors, whose effects are quantified by T
	- 2. The beam reflects off mirror 3  $(M_3)$  and is sent down the declination axis of the telescope, resulting in a rotation in dec,  $R(dec)$
	- 3. The beam reflects off mirror 4  $(M_4)$  and is sent down the polar axis of the telescope, resulting in a rotation in hour angle,  $R(HA)$
	- 4. The light passes through all stationary optical elements past M4, whose effects are quantified by  $M_{rest}$



• Starting with the telescope Mueller matrix T (combined effects of primary and secondary mirrors), the signal after phase correction is:

 $\vec{S}' = \mathsf{MS} = [\mathsf{M}_{\text{rest}} \mathsf{R}(\mathsf{HA}) \mathsf{M}_{4} \mathsf{R}(\mathsf{dec}) \mathsf{M}_{3} \mathsf{T}] \ \vec{S}$ 

- Trace the path of the light:
	- 1.  $\vec{S}$  is altered as the beam passes through the primary and secondary mirrors, whose effects are quantified by T
	- 2. The beam reflects off mirror 3  $(M<sub>3</sub>)$  and is sent down the declination axis of the telescope, resulting in a rotation in dec,  $R(dec)$
	- 3. The beam reflects off mirror 4  $(M_4)$  and is sent down the polar axis of the telescope, resulting in a rotation in hour angle,  $R(HA)$
	- 4. The light passes through all stationary optical elements past M4, whose effects are quantified by  $M_{rest}$
	- 5. The light reaches the camera



Signal after phase correction:

Phase correction  
\n
$$
\vec{S'} = M(HA, dec)\vec{S} = [M_{rest}R(HA)M_4R(dec)M_3\vec{T}] \vec{S}
$$
\nmeasured signal  $\vec{s'}$   
\n
$$
M_{\phi}\vec{S'} = [M_{rest}R(HA)M_4R(dec)M_3] \vec{S}
$$
\n
$$
M(HA, dec)
$$

• Now we can (theoretically) find the "true" Stokes vector  $\vec{S}$  from the measured signal  $\vec{S}^{\prime\prime}$ :

$$
\vec{S} = \mathbf{M}^{-1}(\text{HA}, \text{dec})\mathbf{M}_{\phi}\vec{S}^{\prime\prime}
$$

$$
= [\mathbf{M}_{\text{rest}}\mathbf{R}(\text{HA})\mathbf{M}_4\mathbf{R}(\text{dec})\mathbf{M}_3]^{-1}\mathbf{M}_{\phi}\vec{S}^{\prime\prime}
$$



### Finding calibration matrix  $M_{\text{rest}}$

• We can measure  $M(HA, dec<sub>0</sub>)$  (details to come) over one day, and use this to find  $M_{rest}$ :





NIRIS pipeline Calibration Calibration

### Measuring  $M(HA_i, dec_0)$



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Theoretical  $I, Q, U, V$  signals:

 $\vec{S}_i'$  $\vec{g}_i' = \mathbf{M}(\mathrm{HA}_i,\mathrm{dec}_0)\mathbf{M}_{QWP}\mathbf{M}_{LP}\vec{S}_i$ Unknowns: All 16 elements of  $M(HA_i, dec_0)$ Unknowns:  $\theta_{LP}$  orientation of LP transmission axis  $\theta_{OWP}$  orientation of QWP fast axis  $\delta_{OWP}$  QWP surface retardation

Fit all these parameters to the measured  $I, Q, U, V$  signals (red)

# Simulating  $M(HA, dec)$  from  $M_{rest}$

- Take calibration measurements throughout the day, once per year, to find  $M(HA, dec_0)$
- Find  $M_{rest}$  from  $M(HA, dec_0)$
- Use  $M_{rest}$  to calculate  $M(HA, dec)$  for any combination of HA and dec (i.e. any day of the year)
	- $M(HA, dec) = M_{ref}R(HA)M_{4}R(dec)M_{3}$
- Apply a different calibration matrix for each observation date/time



NIRIS pipeline Calibration Calibration

# 2024/05/19 calibrated data





#### Back to the NIRIS pipeline…

- Surprise! Another crosstalk correction
	- A finer correction for  $Q \leftrightarrow V$  and  $U \leftrightarrow V$  crosstalk
		- Notable residual crosstalk in images with strong magnetic fields (ARs)
	- Choose crosstalk weights that optimize the symmetry (antisymmetry) of the  $Q, U$  (V) Stokes profiles
		- Deviations from symmetric (antisymmetric)  $Q, U, (V)$  profiles are indicative of additional crosstalk
		- Iterative process
- Image processing finished! Results:
	- cal processed images after Mueller matrix crosstalk calibration (QS)
	- cals processed images after iterative crosstalk correction (AR)



### Part III: Inverting NIRIS Data



#### Stokes inversions using NIRIS data

- Straight forward problem:
	- Assumptions:
		- Properties of the solar atmosphere are known
		- Four Stokes profiles are unknown
	- Solve the radiative transfer equation (RTE) directly to find Stokes profiles
- Inverse problem:
	- For our Solar observations:
		- Four Stokes profiles are known (measured)
		- Properties of the solar atmosphere are unknown
	- Requires inversion of the RTE



#### Stokes inversions using NIRIS data

• Radiative transfer equation:

$$
\frac{d\vec{l}}{d\tau_c} = \widetilde{K}(\vec{l} - \vec{S})
$$

- $\vec{l}$  $\vec{I}$  Stokes pseudo-vector:  $\vec{I} \equiv [I, Q, U, V]^T$
- $\tau_c$  optical depth at continuum wavelength
- $-\vec{S}$  source function vector
- $\tilde{K}$  propagation matrix; accounts for:
	- Absorption withdrawal of same amount of energy from all polarization states
	- Pleochromism differential absorption for the polarization states
	- Dispersion transfer among polarization states
- Mapping between two spaces:
	- measurements + instrumental error  $\rightarrow$  physical quantities + uncertainties
- Mapping represents the physics that generates observables from given physical conditions in the object
	- Same observable can yield different results in inversion, depending on assumed underlying physics



### Stokes inversions using NIRIS data

- Model atmosphere parameters:
	- $T$  Temperature
	- $p$  Pressure
	- $v_{LOS}$  Bulk line-of-sight velocity field
	- $\vec{B}$  Magnetic vector field
		- $B$  Strength
		- $\nu$  Inclination with respect to LOS
		- $\varphi$  Azimuth
	- Ad hoc variables
		- $\xi_{mic}$  Micro-turbulence velocity
		- $\xi_{mac}$  Macro-turbulence velocity
		- $f$  Filling factor
- Inversion methods:
	- Local thermodynamic equilibrium (LTE)
	- Weak-field
	- Micro-structured magnetic atmospheres (MISMA)
	- Milne-Eddington (ME)



#### Summary

- NIRIS targets Zeeman-sensitive lines in the NIR to obtain information about Solar magnetic fields and other quantities of interest
- NIRIS uses dual-beam polarimetry to measure Stokes profiles
	- These profiles are used in Stokes inversion methods to obtain magnetic field vector maps
- Instrumental cross-talk needs to be carefully calibrated to yield accurate Stokes inversion results
	- Mueller calculus
	- Iterative fine correction
- The newly-implemented NIRIS cross-talk calibration method works!

